



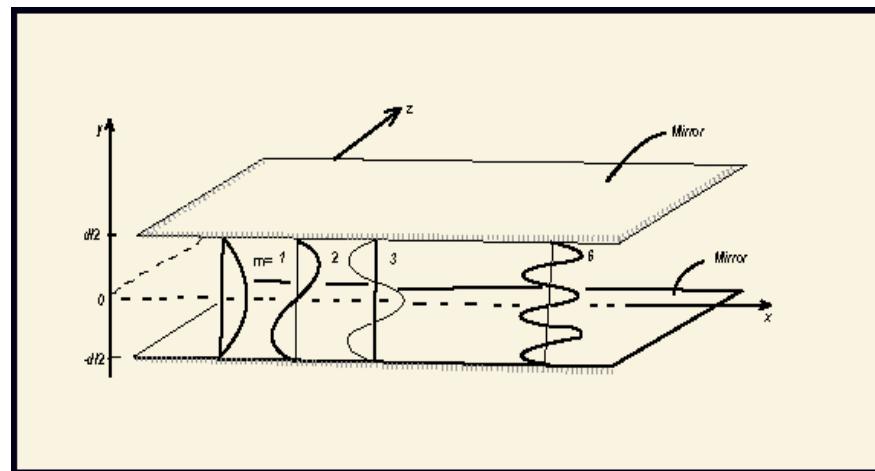
Appendix: the Helmholtz Equation

- Helmholtz Equation
- Finite Difference Method

Guided Wave Optics

Confining Light in 2 out of 3 Dimensions...

- What if we could *trap* a Traveling wave along one/two directions?
 - Standing wave: bounded by a different stiffness (permittivity ϵ)



1D trapping of wave : “**Confinement**”

- What might happen in the trap-free direction (z dimension)?

Guided Wave Optics

Waveguides

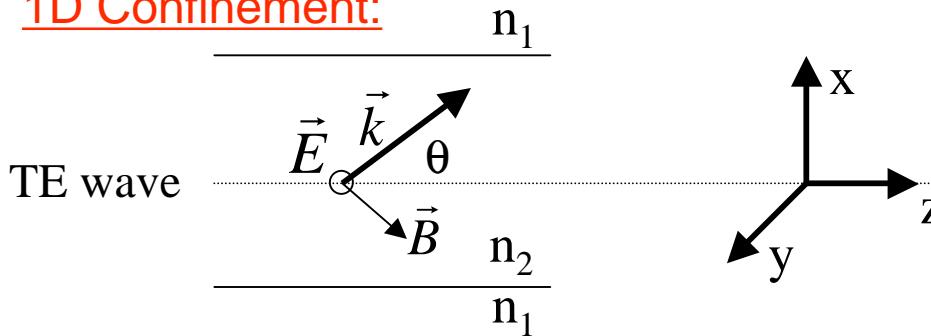
Maxwell's Equations (case: $\rho = 0, J_f = 0$)

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon} = 0 \quad \vec{\nabla} \cdot \vec{B} = 0$$

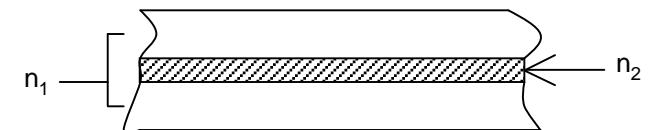
$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \vec{\nabla} \times \vec{B} = \mu \vec{J}_f + \epsilon \mu \frac{\partial \vec{E}}{\partial t} = \epsilon \mu \frac{\partial \vec{E}}{\partial t}$$

- Additional constraint: no propagation in x & y directions
 - Is this even possible?
 - What can we manipulate: n_1, n_2, t_2 (assume $t_1 \rightarrow \infty$)

1D Confinement:



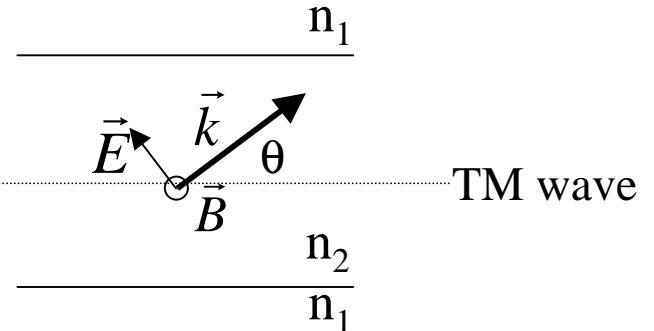
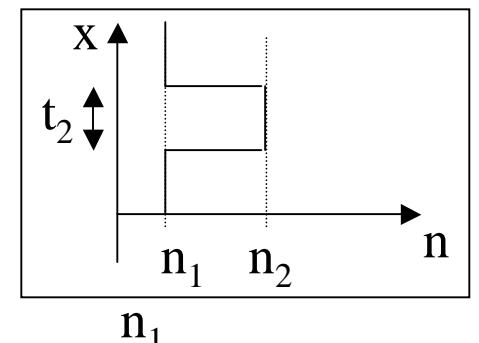
- $E_y \neq 0, E_x = 0, E_z = 0 \Rightarrow E_y = f(x)$
- $B_y = 0, B_x \neq 0, B_z \neq 0 \Rightarrow B_x, B_z = f(x)$



$$\text{Identity: } \vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = \vec{\nabla}(\vec{\nabla} \cdot \vec{E}) - \nabla^2 \vec{E}$$

Traveling
Wave
Equation:

$$\nabla^2 \vec{E} = \epsilon \mu \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{n^2}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$



- $E_y = 0, E_x \neq 0, E_z \neq 0 \Rightarrow E_x, E_z = f(x)$
- $B_y \neq 0, B_x = 0, B_z = 0 \Rightarrow B_y = f(x)$

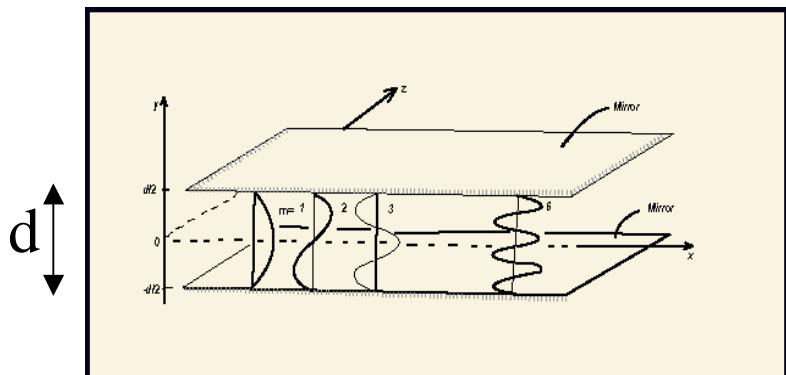
Guided Wave Optics

Waveguide Solution

(TE) Ansatz: $E_y = U(x)e^{-i(\beta z - \omega t)}$

$$\nabla^2 \vec{E} = \frac{n^2}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

Traveling Wave Eqn



Helmholtz Equation:

$$\left[\frac{d^2}{dx^2} + k^2 - \beta^2 \right] U(x) = 0, \quad k = nk_0 = n \frac{\omega}{c}$$

scalar eqn

Standing waves: no more propagation

$$e^{-ikx+i\omega t} + e^{ikx+i\omega t} = 2 \cos(kx)e^{i\omega t}$$

$$U(x) = \begin{cases} A \sin(hx) + B \cos(hx) & |x| < \frac{1}{2}d \\ C \exp(-qx) & x > \frac{1}{2}d \\ D \exp(qx) & x < \frac{1}{2}d \end{cases}$$

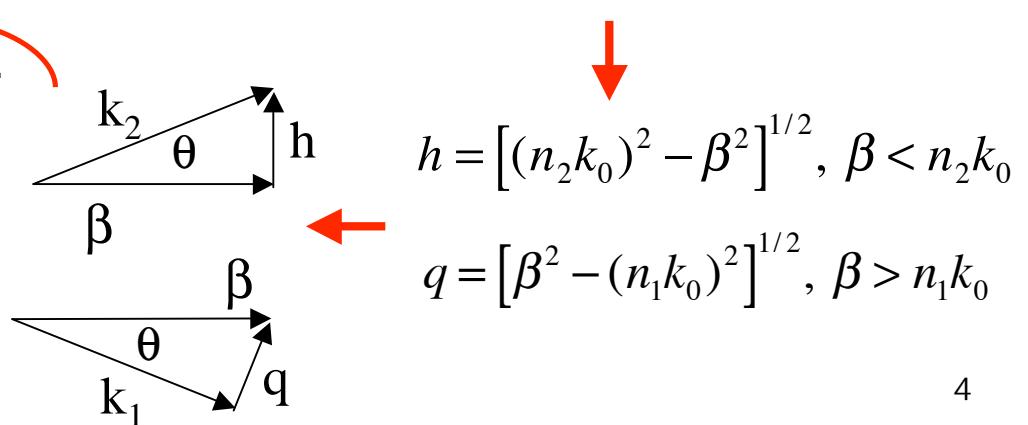
Ray Optic Intuition:

- k_2 bounces around inside wg...
- $k_1 < \beta < k_2$

Effective Index: $\beta = n_{\text{eff}} k_0$

$$\Rightarrow n_1 < n_{\text{eff}} < n_2 \quad (n_{\text{eff}} = n_2 \cos \theta) \\ (\beta = n_2 k_0 \cos \theta)$$

Excerpts: MIT course #3.46



Guided Wave Optics

Waveguide Solution

$$(ii) \vec{E}_{1\parallel} = \vec{E}_{2\parallel} \quad (\text{continuity})$$

$$(iv) \frac{1}{\mu_1} \vec{B}_{1\parallel} = \frac{1}{\mu_2} \vec{B}_{2\parallel} \rightarrow \frac{i}{\omega\mu} \frac{\partial E_x}{\partial y} \Big|_1 = \frac{i}{\omega\mu} \frac{\partial E_x}{\partial y} \Big|_2 \quad (\text{differentiability})$$

4 b'dry conditions:

$$\begin{aligned} A \sin(hd/2) + B \cos(hd/2) &= C \exp(-qd/2) \\ hA \cos(hd/2) - hB \sin(hd/2) &= -qC \exp(-qd/2) \\ -A \sin(hd/2) + B \cos(hd/2) &= D \exp(-qd/2) \\ hA \cos(hd/2) + hB \sin(hd/2) &= qD \exp(-qd/2) \end{aligned}$$

Define Unitless Parameters:

$$u \equiv hd/2, v \equiv qd/2$$

$$B = 0, h \tan(hd/2) = q \quad A = 0, h \cot(hd/2) = -q$$

Odd solns Even solns

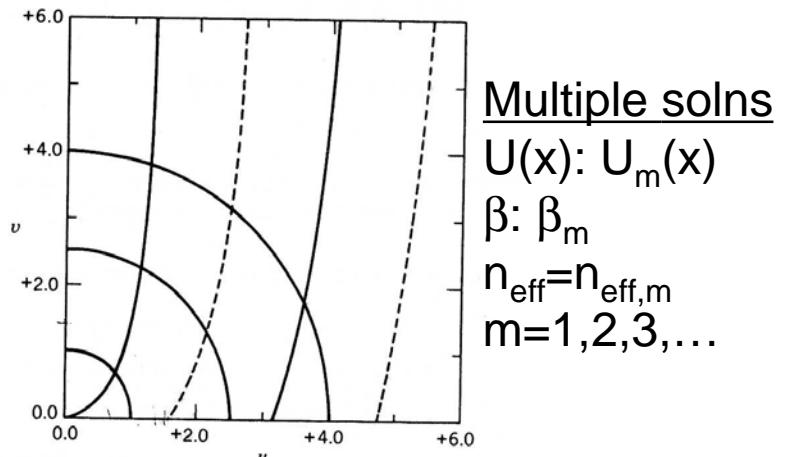
$$\begin{cases} \text{TE: } u \tan(u) = v, u \cot(u) = -v \\ \text{TM: } u \tan(u) = \left(\frac{n_2}{n_1}\right)^2 v, u \cot(u) = -\left(\frac{n_2}{n_1}\right)^2 v \end{cases}$$

$$\text{Identity: } u^2 + v^2 = (n_2^2 - n_1^2) \left(\frac{\omega d}{2c} \right)^2$$

$$= NA^2 \left(\frac{\pi d}{\lambda_0} \right)^2 \equiv V^2$$

"V-parameter"

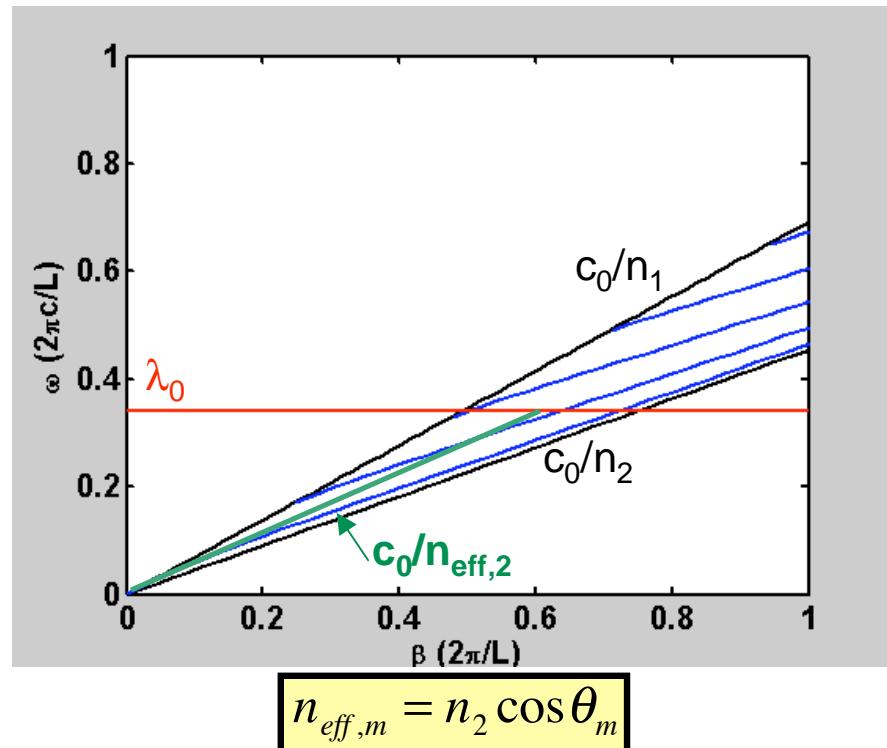
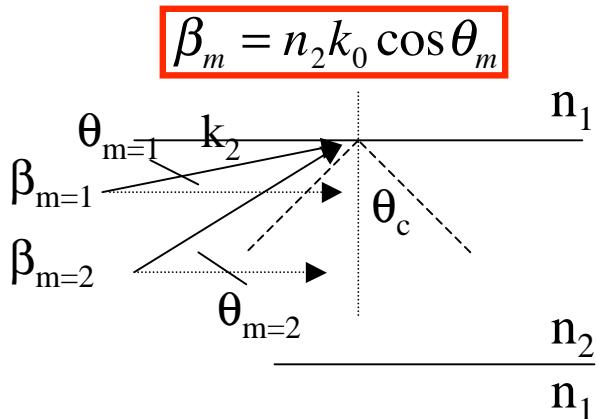
Excerpts: MIT course #3.46



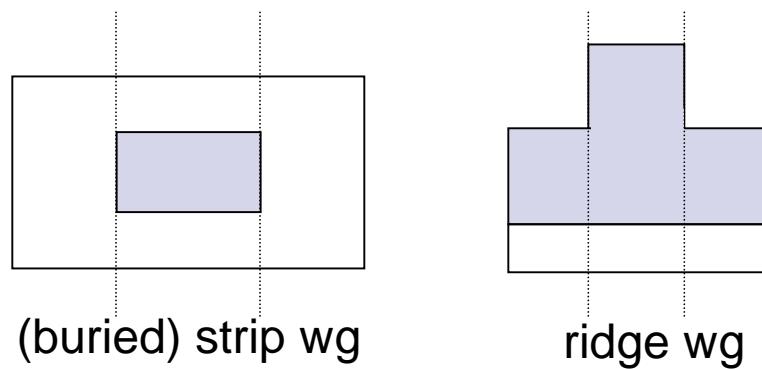
Guided Wave Optics

Waveguide Modes

- Multiple modes ($m=1,2,3,\dots$)
- Single-mode cut-off



- How to deal with 2D confinement?
 - Effective index approach
 - Solve (up to 3) 1-D Eqns



Guided Wave Optics

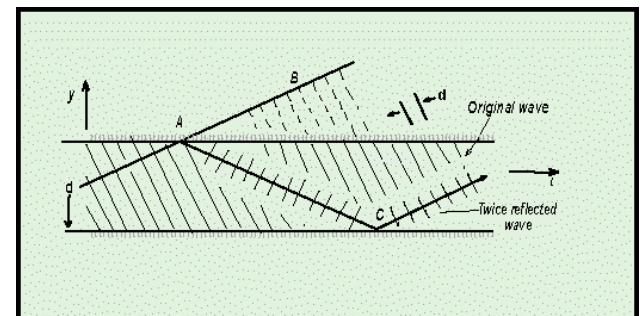
A Ray Optics and Wave Optics picture of Waveguides

$$\beta_m = n_2 k_0 \cos \theta_m$$

- Alternative approach to transcendental eqn:
Total Internal Reflection

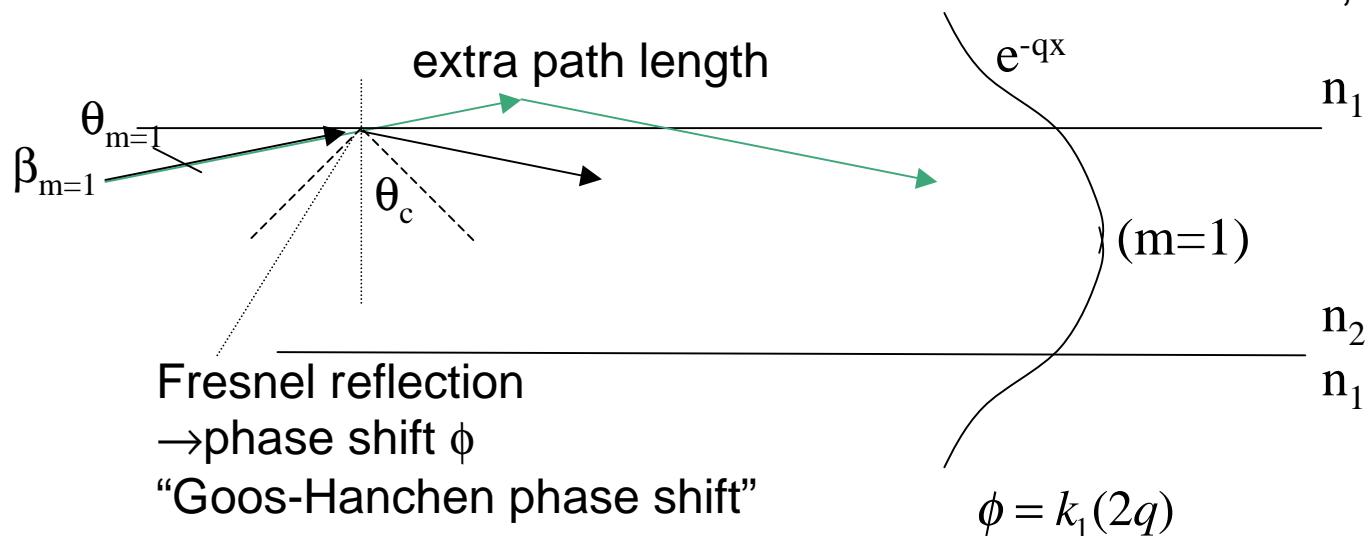
- Ray Optic + Fresnel Eqn: solve for $\cos \theta_m$
(‘Fund. of Phot.’ sec.7.2)
- Bragg’s Law: after two reflections,
experience 2π -phase shift

- Ray Optic visualization



$$m(\lambda_0/n_2) = 2ds \sin \theta_m$$

$$m=1, 2, 3, \dots$$



The Scalar Field Profile

Finite Difference Method

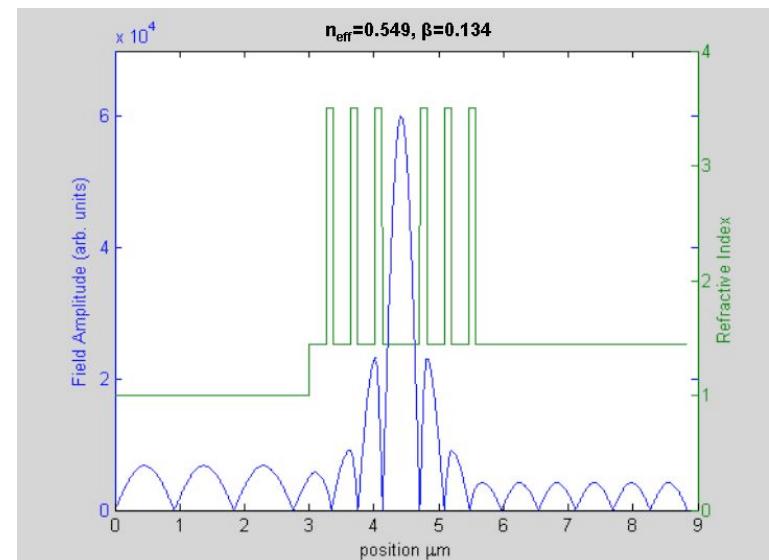
$$\frac{d^2}{dx^2} U_m(y) + k^2 U_m(y) = \beta_m^2 U_m(y) \rightarrow \frac{U_{i+1}}{(\Delta x)^2} + \left[k_0^2 n_i^2 - \frac{2}{(\Delta x)^2} \right] U_i + \frac{U_{i-1}}{(\Delta x)^2} = k_0^2 n_{eff}^2 U_m$$

$$\begin{bmatrix} n_1^2 - 2/(k_0 \Delta x)^2 & 1/(k_0 \Delta x)^2 & 0 & \dots & 0 \\ 1/(k_0 \Delta x)^2 & n_2^2 - 2/(k_0 \Delta x)^2 & 1/(k_0 \Delta x)^2 & 0 & \vdots \\ 0 & 1/(k_0 \Delta x)^2 & \ddots & 1/(k_0 \Delta x)^2 & 0 \\ \vdots & 0 & 1/(k_0 \Delta x)^2 & n_{m-1}^2 - 2/(k_0 \Delta x)^2 & 1/(k_0 \Delta x)^2 \\ 0 & \dots & 0 & 1/(k_0 \Delta x)^2 & n_m^2 - 2/(k_0 \Delta x)^2 \end{bmatrix} \begin{pmatrix} U_1 \\ U_2 \\ \vdots \\ U_{m-1} \\ U_m \end{pmatrix} = n_{eff}^2 \begin{pmatrix} U_1 \\ U_2 \\ \vdots \\ U_{m-1} \\ U_m \end{pmatrix}$$

$$U_{1-1} = 0, U_{m+1} = 0$$

- Discretize x-axis
 - Helmholtz: differential → difference eqn
 - $U(x) \rightarrow U_i$
 - $U_i = f(U_{i-1}, U_{i+1})$

(Applicable to TE-mode E-field profile)



Δx